

ASTROCHEMISTRY LECTURE COURSE PROBLEMS

①

Q1 a) Energy of first transition in Lyman, Balmer, Paschen, Brackett, - Pfund series of atomic hydrogen

$$E = hcR\left(\frac{1}{n_1^2} - \frac{1}{n_2^2}\right) \quad \text{Rydberg equation}$$

$$R = 109737 \text{ cm}^{-1}$$

Lyman $n_1=1$
 $n_2=2 \Rightarrow E = hcR\left(1 - \frac{1}{4}\right) = 1.635 \times 10^{-18} \text{ J}$
 $\lambda = \frac{hc}{E} = 121.5 \text{ nm}$

Balmer $n_1=2$
 $n_2=3 \Rightarrow E = hcR\left(\frac{1}{4} - \frac{1}{9}\right) = 3.028 \times 10^{-19} \text{ J}$
 $\lambda = 656 \text{ nm}$

Paschen $n_1=3$
 $n_2=4 \Rightarrow E = hcR\left(\frac{1}{9} - \frac{1}{16}\right) = 1.060 \times 10^{-19} \text{ J}$
 $\lambda = 1.875 \mu\text{m}$

Brackett $n_1=4$
 $n_2=5 \Rightarrow E = hcR\left(\frac{1}{16} - \frac{1}{25}\right) = 4.905 \times 10^{-20} \text{ J}$
 $\lambda = 4.05 \mu\text{m}$

Pfund $n_1=5$
 $n_2=6 \Rightarrow E = hcR\left(\frac{1}{25} - \frac{1}{36}\right) = 2.664 \times 10^{-20} \text{ J}$
 $\lambda = 7.456 \mu\text{m}$

b) Atmospheric windows span 300-900 nm, 1-5 μm , and 8-20 μm , so the first transitions in the Balmer, Paschen, & Brackett series could be observed using a ground-based telescope.

Q2 a) For the H_α $n=2 \rightarrow 3$ transition, we have

$$E = hcR\left(\frac{1}{4} - \frac{1}{9}\right) = 3.0275597 \times 10^{-19} \text{ J}$$

$$\lambda = 656.114 \text{ nm}$$

b) The Doppler shift is $\Delta\lambda = \left(\frac{v_{\text{source}}}{c}\right)\lambda$

The rotation speed is 43000 km hr^{-1}
 $= 43000 (1000 \text{ m}) (3600 \text{ s})^{-1}$
 $= 11944 \text{ ms}^{-1}$

The Doppler shifts of points on the left and right of Jupiter, moving towards and away from Earth at this speed, are ②

$$\begin{aligned}\Delta\lambda &= \frac{v_{\text{source}}}{c} \lambda \\ &= \pm \frac{11944}{c} \cdot 686.114 \\ &= \pm 0.026 \text{ nm}\end{aligned}$$

A point observed in the centre of the planet as viewed from Earth has zero velocity component along the line of sight, and $\therefore$ no Doppler shift due to rotation (there may be a small Doppler shift due to the relative motion of the two planets in their orbits)

Q3. From the Boltzmann distribution, we expect

$$\frac{N_1}{N_0} = \frac{g_1 e^{-E_1/kT}}{g_0 e^{-E_0/kT}} = 3 e^{-(E_1-E_0)/kT}$$

The transition energy is (ignoring centrifugal distortion)
 $E_1 - E_0 = 2B(J+1) = 2B$

$$B = \frac{h}{8\pi^2 c I}, \text{ and for a diatomic, } I = \mu r^2.$$

The reduced mass of CO is

$$\begin{aligned}\mu &= \frac{m m_0}{m + m_0} = \frac{12 \times 16}{12 + 16} = 6.857 \text{ g mol}^{-1} \\ &= 1.139 \times 10^{-26} \text{ kg}\end{aligned}$$

The bond length is

$$r = 112.8 \text{ pm} = 112.8 \times 10^{-12} \text{ m}$$

$$\begin{aligned}\Rightarrow I &= (1.139 \times 10^{-26}) (112.8 \times 10^{-12})^2 \\ &= 1.449 \times 10^{-46} \text{ kg m}^2\end{aligned}$$

$$\Rightarrow B = \frac{h}{8\pi^2 c I} = 1.932 \text{ cm}^{-1}$$

The transition E is $\therefore$

$$E_1 - E_0 = 2B = 3.864 \text{ cm}^{-1} = 7.676 \times 10^{-23} \text{ J}$$

(3)

Rearranging the equation for $\frac{n_1}{n_0}$ gives

$$\ln\left(\frac{1}{3} \frac{n_1}{n_0}\right) = -\frac{(E_1 - E_0)}{kT}$$

$$\begin{aligned} T &= -\frac{(E_1 - E_0)}{k \ln\left(\frac{1}{3} \frac{n_1}{n_0}\right)} \\ &= -\frac{7.676 \times 10^{-23}}{k \ln\left(\frac{1.10}{3}\right)} \\ &= 8.54 \text{ K} \end{aligned}$$

Q4. From Q 2(a), the transition energy for the $H\alpha$ line is $3.0276 \times 10^{-19} \text{ J}$. The ratio of populations for $n=2$ and $n=3$ at 11000 K is $\therefore$

$$\begin{aligned} \frac{n_3}{n_2} &= \frac{g_3}{g_2} e^{-\Delta E/kT} \\ &= \frac{9}{4} \exp\left(\frac{-3.0276 \times 10^{-19}}{11000 k}\right) \\ &= 0.31 \end{aligned}$$

Q5 a) See lecture notes

b) Bond dissociation occurs when the thermal energy kT exceeds the bond dissociation energy i.e.

$$T > \frac{E_{\text{bond}}}{k}$$

$$\text{(i) N-N} \quad T > \frac{943000}{8.314} = 1.134 \times 10^5 \text{ K}$$

$$\text{(ii) C=O} \quad T > \frac{1075000}{8.314} = 1.293 \times 10^5 \text{ K}$$

$$\text{(iii) C=C} \quad T > \frac{612000}{8.314} = 73,610 \text{ K}$$

$$\text{(iv) C-C} \quad T > \frac{348000}{8.314} = 41857 \text{ K}$$

$$\text{(v) C-H} \quad T > \frac{415000}{8.314} = 49916 \text{ K}$$

(4)

(c) Similarly, atoms are ionized when

$$T > \frac{I.P.}{k}$$

$$(i) \text{ H } T > \frac{13.6 \text{ eV}}{k} = 157809 \text{ K}$$

$$(ii) \text{ He } T > \frac{24.59 \text{ eV}}{k} = 285333 \text{ K}$$

$$(iii) \text{ C } T > \frac{11.26 \text{ eV}}{k} = 130657 \text{ K}$$

$$(iv) \text{ He}^+ T > \frac{54.51 \text{ eV}}{k} = 632513 \text{ K}$$

Q6. For a density of $1000 \text{ cm}^{-3} = 1 \times 10^9 \text{ m}^{-3}$, a mass of $2 \times 10^{30} \text{ kg}$ would take up a volume of

$$V = \frac{m}{\rho}$$

$$= \frac{m}{m_H n}$$

$\swarrow$ mass of H atom $\nwarrow$ number density

NB: Assume the cloud is made up of atomic hydrogen

$$= \frac{2 \times 10^{30}}{(1.66 \times 10^{-27})(1 \times 10^9)}$$

$$= 1.204 \times 10^{48} \text{ m}^3$$

Assuming that the cloud is spherical, we have

$$V = \frac{4}{3} \pi r^3$$

$$r = \left(\frac{3V}{4\pi} \right)^{\frac{1}{3}}$$

$$= \left(\frac{3 \cdot 1.204 \times 10^{48}}{4\pi} \right)^{\frac{1}{3}}$$

$$= 6.600 \times 10^{15} \text{ m}$$

The diameter is $\therefore$

$$d = 2(6.600 \times 10^{15} \text{ m})$$

$$= 1.32 \times 10^{16} \text{ m}$$

$$= 1.4 \text{ light years}$$

Q7. The collision frequency is given by

$$Z = \frac{\sigma_c}{2} \left(\frac{8kT}{\pi m} \right)^{\frac{1}{2}} n^2$$

For H_2 , we have

$$\begin{aligned}\sigma_c &= 0.27 \text{ nm}^2 \\ &= 2.7 \times 10^{-19} \text{ m}^2\end{aligned}$$

$$m = 8.303 \times 10^{-28} \text{ kg}$$

(a) Using the ideal gas law, the number density at atmospheric pressure and room T

$$n = \frac{N}{V} = \frac{P}{kT} = \frac{101300}{298k} = 2.45 \times 10^{25} \text{ m}^{-3}$$

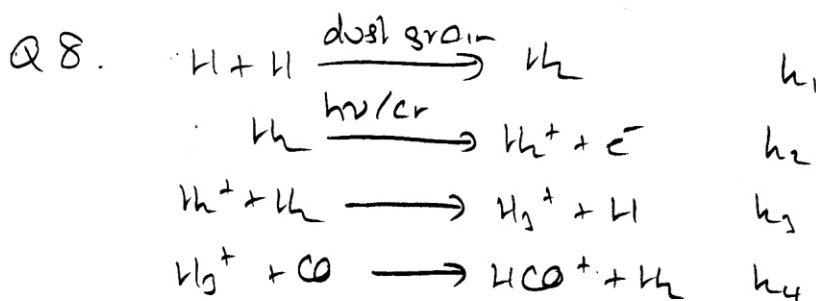
$$\begin{aligned}\Rightarrow Z &= \left(\frac{2.7 \times 10^{-19}}{2} \right) \left(\frac{8k \cdot 300}{8.303 \times 10^{-28} \pi} \right)^{\frac{1}{2}} (2.45 \times 10^{25})^2 \\ &= 2.89 \times 10^{35} \text{ s}^{-1} \text{ m}^{-3}\end{aligned}$$

(b) when $n = 10^5 \text{ cm}^{-3} = 1 \times 10^{11} \text{ m}^{-3}$, and $T = 40 \text{ K}$, we have

$$\begin{aligned}Z &= \left(\frac{2.7 \times 10^{-19}}{2} \right) \left(\frac{8k \cdot 40}{8.303 \times 10^{-28} \pi} \right)^{\frac{1}{2}} (1 \times 10^{11})^2 \\ &= 1.76 \times 10^6 \text{ m}^{-3} \text{ s}^{-1}\end{aligned}$$

(c) when $n = 10 \text{ cm}^{-3} = 1 \times 10^7 \text{ m}^{-3}$, and $T = 10 \text{ K}$, we have

$$\begin{aligned}Z &= \left(\frac{2.7 \times 10^{-19}}{2} \right) \left(\frac{8k \cdot 10}{8.303 \times 10^{-28} \pi} \right)^{\frac{1}{2}} (1 \times 10^7)^2 \\ &= 0.0088 \text{ m}^{-3} \text{ s}^{-1} \\ &= 32 \text{ m}^{-3} \text{ hr}^{-1}\end{aligned}$$



(6)

a) The reaction steps are:

1. Surface catalysed reaction of 2 H atoms on a dust grain to form H_2
2. Photoionisation or electron-impact ionization (an e^- or photon knocks an e^- out of H_2 to make H_2^+)
3. Ion-molecule reaction (H atom abstraction from H_2 by H_2^+) to form H_3^+ .
4. Ion-molecule reaction (proton transfer from H_3^+ to CO) to form HCO^+ .

b) Apply SSA to $[H_2^+]$ and $[H_3^+]$

$$\frac{d[H_2^+]}{dt} = 0 = k_2[H_2] - k_3[H_2^+][H_2]$$

$$\Rightarrow [H_2^+] = \frac{k_2}{k_3}$$

$$\frac{d[H_3^+]}{dt} = 0 = k_3[H_2^+][H_2] - k_4[H_3^+][CO]$$

$$[H_3^+] = \frac{k_3[H_2^+][H_2]}{k_4[CO]}$$

$$= \frac{k_2[H_2]}{k_4[CO]} \quad \text{since } [H_2^+] = \frac{k_2}{k_3}$$

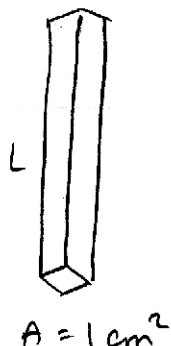
$$\therefore \frac{d[HCO^+]}{dt} = k_4[H_3^+][CO]$$

$$= k_4 \cdot \frac{k_2[H_2]}{k_4[CO]} \cdot [CO]$$

$$= k_2[H_2]$$

Ionisation of H_2 is the rate limiting step in the mechanism.

c)



The no. of molecules in a column of cross section 1 cm^2 and length L is 3×10^{24} .

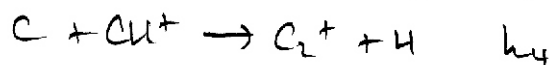
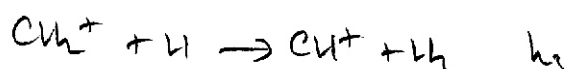
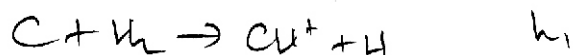
The length of the column is $\therefore$

$$L = \frac{3 \times 10^{24}}{2 \times 10^7} = 1.5 \times 10^{20} \text{ cm} = 1.5 \times 10^{18} \text{ m}$$

This is about 160 light years

(1 light year = $9.46 \times 10^{15} \text{ m}$).

Q9

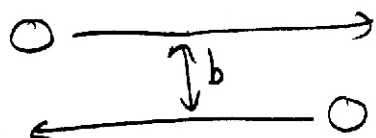


a) Treating the CH^+ ion as a reactive intermediate, we have

$$\frac{d[\text{CH}^+]}{dt} = 0 = k_1[\text{C}][\text{H}_2] - k_2[\text{CH}^+][\text{N}] + k_3[\text{CH}_2^+][\text{H}] - k_4[\text{C}][\text{CH}^+]$$

$$[\text{CH}^+] = \frac{k_1[\text{C}][\text{H}_2]}{k_2[\text{N}] - k_3[\text{H}] + k_4[\text{C}]}$$

b)



The impact parameter is the distance by which two particles would miss each other in the absence of attractive or repulsive interactions. It gives a

measure of how "direct" a collision is. i.e. very small b corresponds to a "head on" collision, larger b corresponds to a "glancing" collision, and much larger b will eventually lead to no collision. The reaction cross section can be defined in terms of the maximum b leading to reaction i.e. $\sigma = \pi b_{\text{max}}^2$

c) The orbital angular momentum is

$$L = \mu v b$$

Taking v as the mean relative velocity,

$$v = \left(\frac{8kT}{\pi \mu} \right)^{\frac{1}{2}}$$

we have

$$L = \mu \left(\frac{8kT}{\pi \mu} \right)^{\frac{1}{2}} b$$

i.e. for a given b and T ,

$$L \propto \mu^{\frac{1}{2}}$$

$$= \left(\frac{8kT\mu}{\pi} \right)^{\frac{1}{2}} b$$

For the four reactions, we have

$$1. \mu(\text{CC}^+ + \text{H}) = \frac{12 \times 2}{12 + 2} = 1.714 \text{ g mol}^{-1} = 2.847 \times 10^{-27} \text{ kg}$$

$$L = \left(\frac{8 \times 6.50 \times 2.847 \times 10^{-27}}{\pi} \right)^{\frac{1}{2}} b$$

$$= (3.94 \times 10^{-24}) b$$

$$2. \mu(\text{CH}^+ + \text{N}) = \frac{13 \times 14}{13 + 14} = 6.74 \text{ g mol}^{-1} = 1.119 \times 10^{-26} \text{ kg}$$

$$L = \left(\frac{8 \times 6.50 \times 1.119 \times 10^{-26}}{\pi} \right)^{\frac{1}{2}} b$$

$$= (4.44 \times 10^{-24}) b$$

$$3. \mu(\text{CCu}^+ + \text{H}) = \frac{14 \times 1}{14 + 1} = 0.933 \text{ g mol}^{-1} = 1.550 \times 10^{-27} \text{ kg}$$

$$L = \left(\frac{8 \times 6.50 \times 1.550 \times 10^{-27}}{\pi} \right)^{\frac{1}{2}} b$$

$$= (1.65 \times 10^{-24}) b$$

$$4. \mu(\text{C} + \text{CH}^+) = \frac{12 \times 13}{12 + 13} = 6.24 \text{ g mol}^{-1} = 1.036 \times 10^{-26} \text{ kg}$$

$$L = \left(\frac{8 \times 6.50 \times 1.036 \times 10^{-26}}{\pi} \right)^{\frac{1}{2}} b$$

$$= (4.27 \times 10^{-24}) b$$

The orbital angular momentum is highest for $Cl^{+} + N$

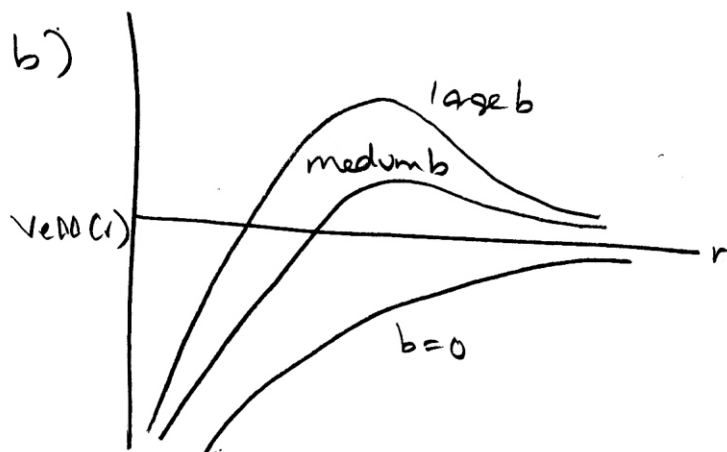
- d) The orbital angular momentum leads to a centrifugal barrier $\frac{L^2}{2I}$ on the PES. The barrier height increases with b and with the reduced mass of the collision partners, so the second step is likely to be the rate determining step.

Q10.
$$V_{eff} = -\frac{\alpha q^2}{8\pi\epsilon_0 r^4} - \frac{M_0 q}{4\pi\epsilon_0 r^2} + \frac{M_{rel}^2 b^2}{2r^2}$$

- a) ion-induced dipole term:
The ionic charge induces a dipole in the neutral, leading to an attractive interaction

ion-dipole term:
The ionic charge interacts with a permanent dipole on the neutral molecule.

centrifugal barrier term:
Angular momentum associated with the orbital motion of the reactants must be conserved. The kinetic E associated with this motion is not available for reaction & is often treated in terms of a effective barrier on the PES.



- c) To find the maximum, we differentiate $V_{eff}(r)$ w.r.t. r , set the result equal to zero, and solve for r to find r_0 (10)

$$\frac{dV_{eff}}{dr} = 0 = \frac{d}{dr} \left[-\frac{\alpha q^2}{8\pi\epsilon_0} r^{-4} - \frac{m_0 q}{4\pi\epsilon_0} r^{-2} + \frac{m v_{rel}^2 b^2}{2} r^{-2} \right]$$

$$= \frac{4\alpha q^2}{8\pi\epsilon_0 r^5} + \frac{2m_0 q}{4\pi\epsilon_0 r^3} - \frac{2m v_{rel}^2 b^2}{2 r^3}$$

$$= \frac{\alpha q^2}{r^5} + m_0 q - 2\pi\epsilon_0 m v_{rel}^2 b^2$$

$$r = \left(\frac{\alpha q^2}{2\pi\epsilon_0 m v_{rel}^2 b^2 - m_0 q} \right)^{\frac{1}{2}}$$

- d) As the impact parameter b increases, the height of the centrifugal barrier increases (see figure in b)). At some value of b the barrier will exceed the collision energy E_{coll} , we can find this maximum impact parameter by setting the barrier height equal to E_{coll} , solving for b , i.e.

$$V_{eff}(r_0) = E_{coll}$$

The reaction cross section is then

$$\sigma_r(E_{coll}) = \pi b_{max}^2$$

- e) From simple collision theory, the rate constant for a given relative velocity is

$$k(v_{rel}) = \sigma_r(v_{rel}) v_{rel}$$

The thermal rate constant is found by integrating $k(v_{rel})$ over the Maxwell-Boltzmann velocity distribution at temperature T , $f(v_{rel})$.

$$k(T) = \int \frac{\sigma_r(v_{rel}) v_{rel}}{k(v_{rel})} f(v_{rel}) dv_{rel}$$